

Rendering: Materials 1

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- Adding refractions
 - Snell's Law
 - Fresnel Reflectance
 - Specular BTDF

- Important concepts
 - Chromatic Aberration
 - Heckbert Notation
 - Caustics



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- ***Physical (wave) optics:***
 - Derived using a detailed model of light
 - Treating it as wave and computing solutions to Maxwell's equations
 - Computationally expensive, usually not appreciably more accurate

- ***Geometric optics:***
 - Requires surface's low-level scattering and geometric properties
 - Closed-form reflection models derived from these properties
 - More tractable, complex wave effects like polarization are ignored



$$L_e(x, v) = E(x, v) + \int_{\Omega} f_r(x, \omega \rightarrow v) L_i(x, \omega) \cos(\theta_x) d\omega$$

- Bidirectional Scattering Distribution Function (BSDF)
- Describes the light transport properties of the material
- So far, we only implemented simple diffuse surface reflections
- Can model reflections, refractions, volumetric scattering...



- Considers only the **reflection** of incoming light onto a surface
 - The BRDF is a limited instance of the full BSDF (e.g., no transparency)
 - Good for starting out, complex materials need full BSDF
 - More on that in another lecture
- A BRDF function $f_r(x, \omega_i \rightarrow \omega_o)$ with input directions ω_i, ω_o
 - uses convention: ω_i and ω_o are assumed to point away from x
 - How much irradiance from ω_i is reflected as radiance to ω_o at x ?



- “How much irradiance from ω_i is reflected as radiance to ω_o at x ?”

- $f_r(x, \omega_i \rightarrow \omega_o) = \frac{dL_i(x, \omega_o)}{dE_i(x, \omega_i)} = \frac{dL_i(x, \omega_o)}{L_i(x, \omega_i) \cos\theta(\omega_i) d\omega_i}$

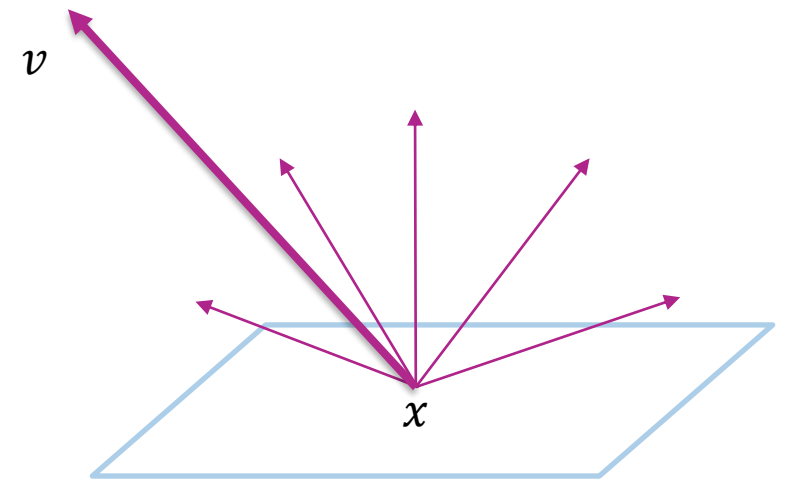
$$L_e(x, v) = E(x, v) + \int_{\Omega} f_r(x, \omega \rightarrow v) L_i(x, \omega) \cos(\theta_x) d\omega$$

- **Helmholtz reciprocity:** $f_r(x, \omega_i \rightarrow \omega_o) = f_r(x, \omega_o \rightarrow \omega_i)$

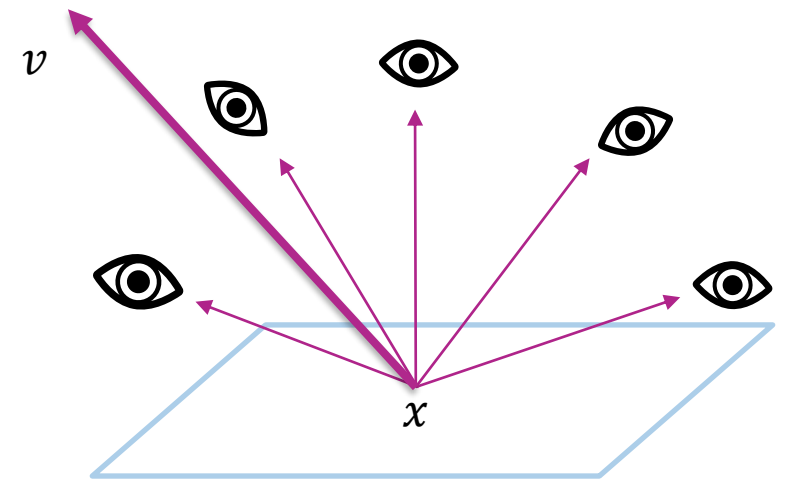
- **Conserves energy:** $\int_{\Omega} f_r(x, \omega \rightarrow v) \cos\theta d\omega \leq 1 \forall v$



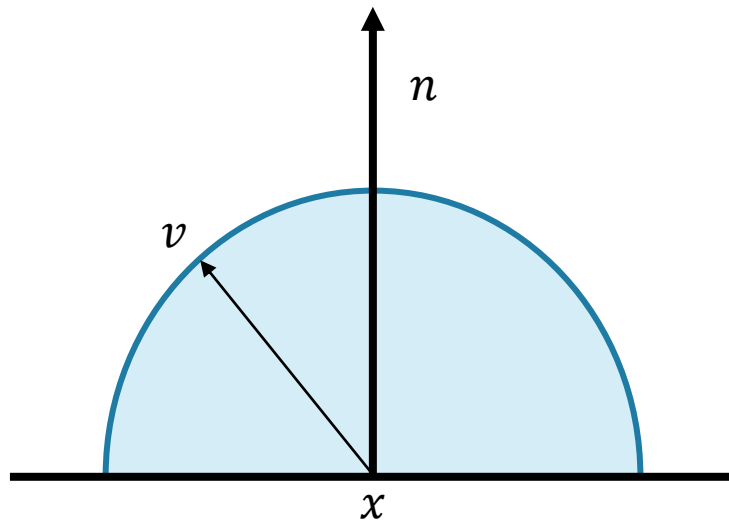
- Why must the BRDF f_r fulfill $\int_{\Omega} f_r(x, \omega \rightarrow v) \cos\theta(\omega) d\omega \leq 1$?
- Intuitive interpretation with **reciprocity**: Shine a laser light along $-v$ onto x . We must have $\int_{\Omega} f_r(x, v \rightarrow \omega) \cos\theta(\omega) d\omega \leq 1$
- If we find a direction v for which this is not true, it means we would reflect more light than is coming in



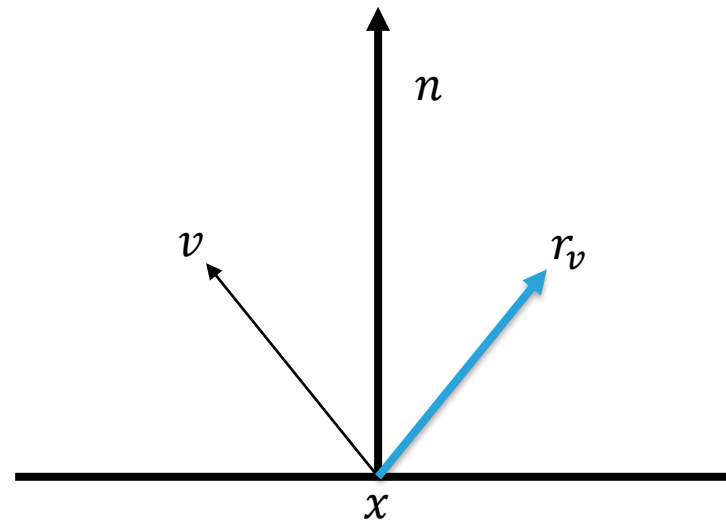
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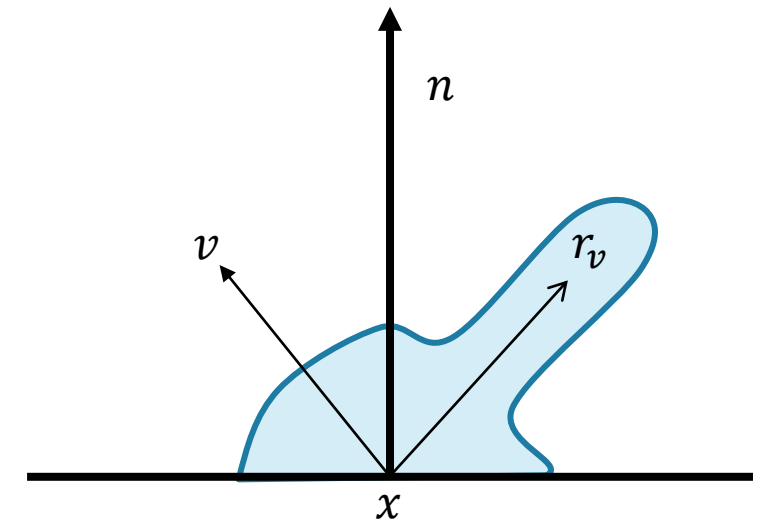
- We usually distinguish three basic BRDF types
 - Perfectly diffuse (light is scattered equally in/from all directions)
 - Perfectly specular (light is reflected in/from exactly one direction)
 - Glossy (mixture of the other two, stronger reflectance around r_v)



Diffuse



Specular



Glossy



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- Before, we considered the BRDF value and sampling of ω separately
- For implementation, it makes a lot of sense to combine them
 - $f_r(x, \omega \rightarrow v)$ depends only on x , v and next ray direction ω
 - Rendering equation: we can't predict L_i , but $f_r(x, \omega \rightarrow v)$ and $\cos \theta$
 - Our renderings will converge faster if the distribution of ω actually matches the shape of $f_r(x, \omega \rightarrow v) \cos \theta$ (**importance sampling!**)
 - If we put the BRDF in charge of choosing our ω , we can make it sample a distribution that directly matches $f_r(x, \omega \rightarrow v) \cos \theta$
 - This actually makes things cleaner in code

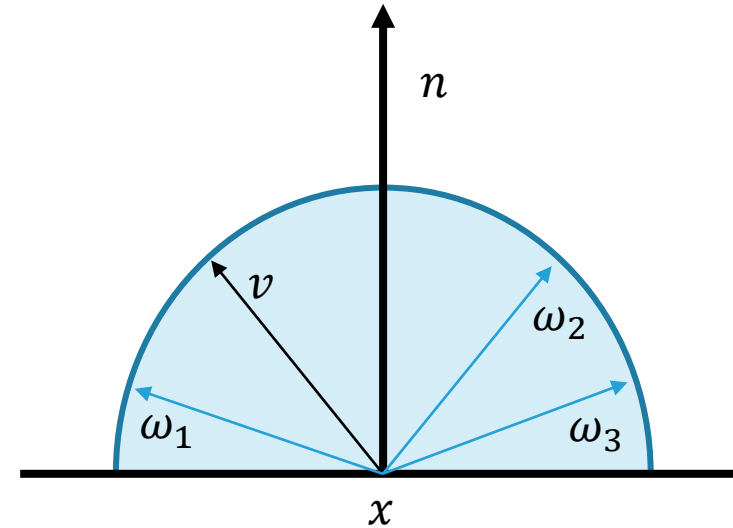


- Diffuse materials reflect same amount of light in/from all directions

- $f_r(x, \omega \rightarrow v) = \frac{\rho}{\pi} \forall v, \omega \angle n < \frac{\pi}{2}$

- ρ = amount of reflected light

- $\rho \leq 1$ in r, g, b



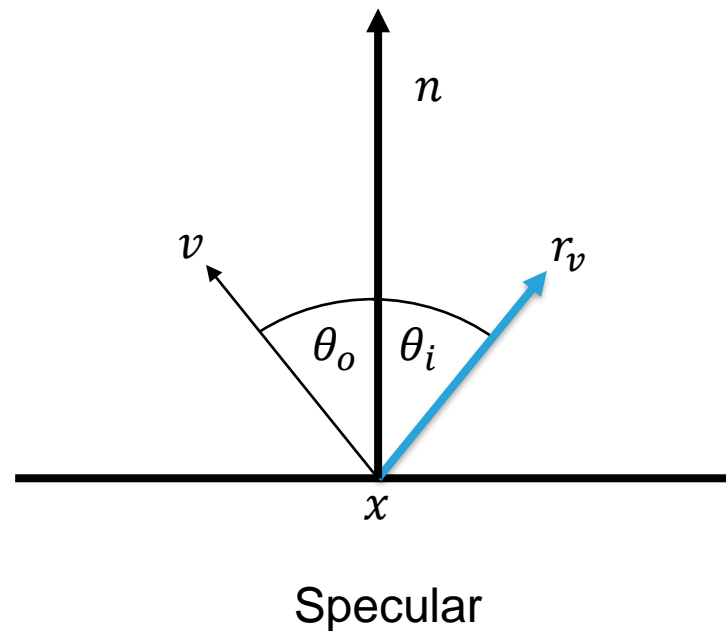
- Importance sampling $f_r(x, \omega \rightarrow v) \cos \theta \rightarrow$ use $p(\omega) \propto \frac{\rho \cos \theta}{\pi}$

- Making it a valid PDF leads to $p(\omega) = \frac{\cos \theta}{\pi}$

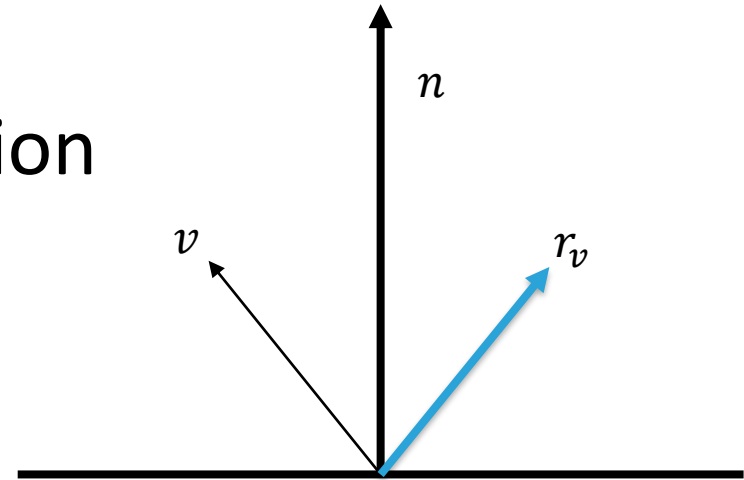
- From previous exercise: it's cosine-weighted hemisphere sampling!



- The angle of exiting light θ_o is the same as the angle of incidence θ_i
- Incoming light is only transported in a single direction



- For purely specular BRDFs (a perfect mirror surface), irradiance from the perfect mirror direction r_v is completely reflected to v
- Irradiance coming from any other direction does not reflect at all towards v
- $f_r(x, \omega \rightarrow v) > 0 \Leftrightarrow \omega = r_v$
- Problem: if we pick the next direction ω randomly as before, the chances of ever hitting r_v by accident are infinitely small!



- Model specular reflection with the Dirac delta function
- Delta function $\delta(x)$ is defined to be 0 everywhere except at $x = 0$
- Use a shifted version $\delta_v(\omega)$ that is 0 everywhere except at $\omega = r_v$
- Per definition, $\int_{\Omega} \delta_v(\omega) d\omega = 1$ to obtain a valid PDF for sampling
- Ponder this for a moment: what value does $\delta_v(r_v)$ have?



- Full energy preservation: $\int_{\Omega} f_r(x, \omega \rightarrow v) L_i \cos_{\theta}(\omega) d\omega = L_{r_v}$
- If we integrate using $f_r(x, \omega \rightarrow v) = \delta_v(\omega)$, we get $L_{r_v} \cos_{\theta}(r_v)$
- We lost some light! We compensate: $f_r(x, \omega \rightarrow v) = \frac{\delta_v(\omega)}{\cos_{\theta}(r_v)}$
- If we consider the properties of the Dirac delta function, we can try to derive the same methods that we used before for diffuse BRDFs



- **sample(v)**: mirror v about n (invert v_x, v_y in *local space*) and return

- **evaluate(a, b)**: 0 if $b \neq r_a$, else return $\frac{\delta_a(r_a)}{\cos\theta(r_a)} = \frac{\infty}{\cos\theta(r_a)}$

- **Problem**: How to calculate anything reasonable with ∞ ?

- **Problem**: we are comparing two vectors with floats (Stability?)

- **pdf(ω)**: 0 if $\omega \neq r_v$, else: $\delta_v(r_v) = \infty$

- But, if $\omega = r_v$, **evaluate(v, ω) / pdf(ω)** = $\frac{\delta_v(\omega)}{\delta_v(\omega)\cos\theta(r_v)} = \frac{1}{\cos\theta(r_v)}$



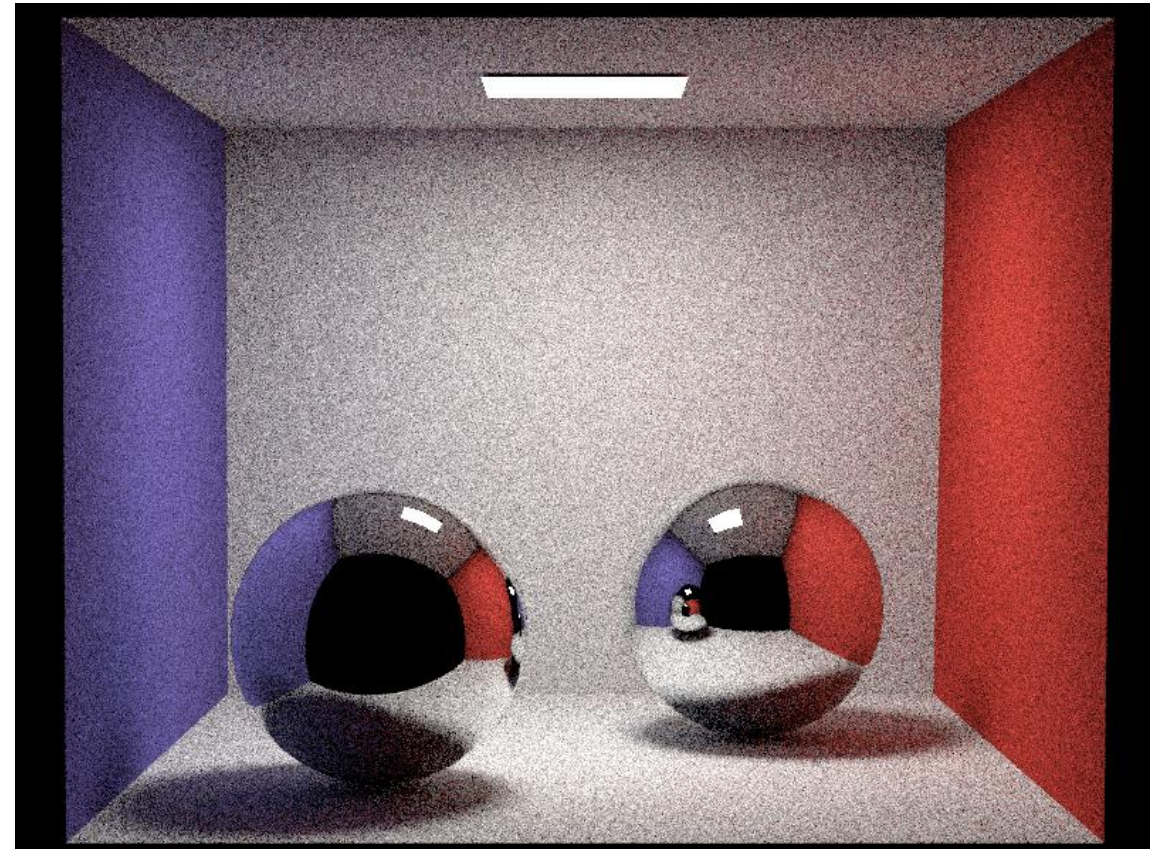
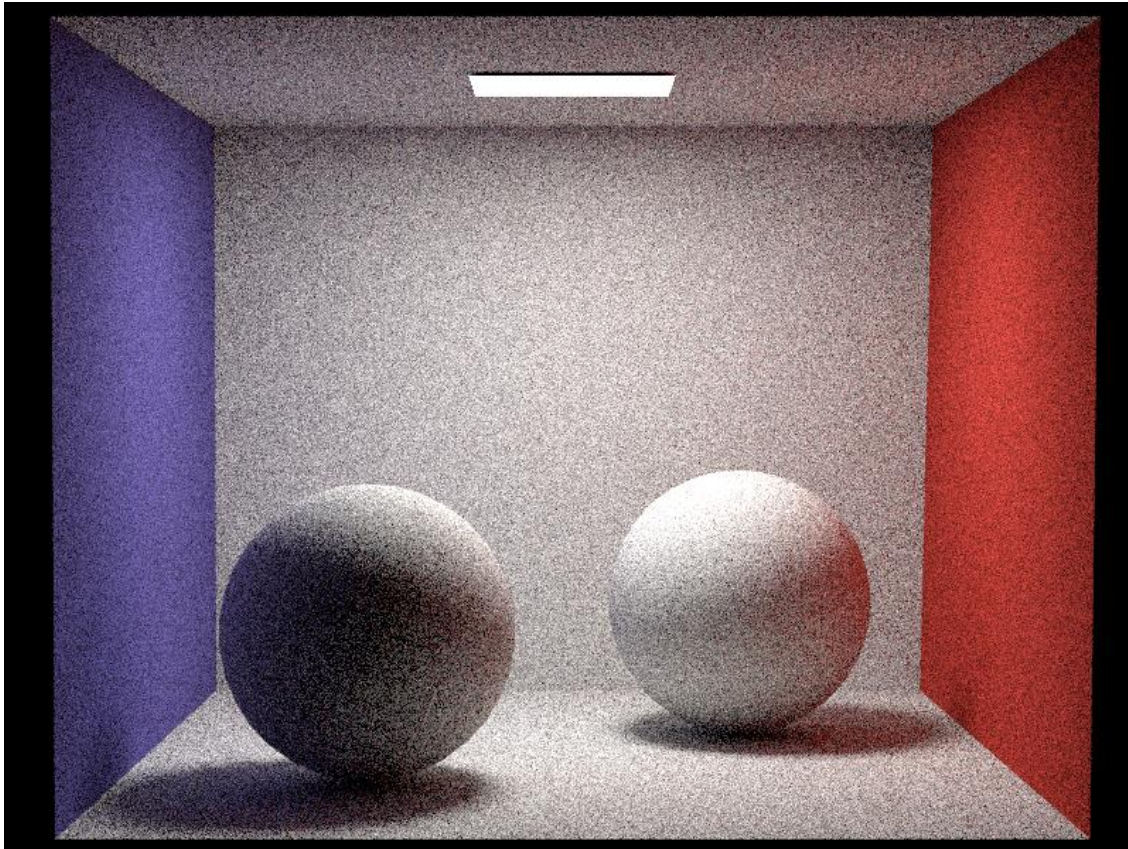
- **sample(v)**: mirror v about n (invert v_x, v_y in *local space*)
 - Return r_v as generated sample direction
 - Return multiplier for L_i as 1 (full radiance passed on)

- No other function except **sample** should be able to just *guess* r_v

- **evaluate(a, b)**: always return 0

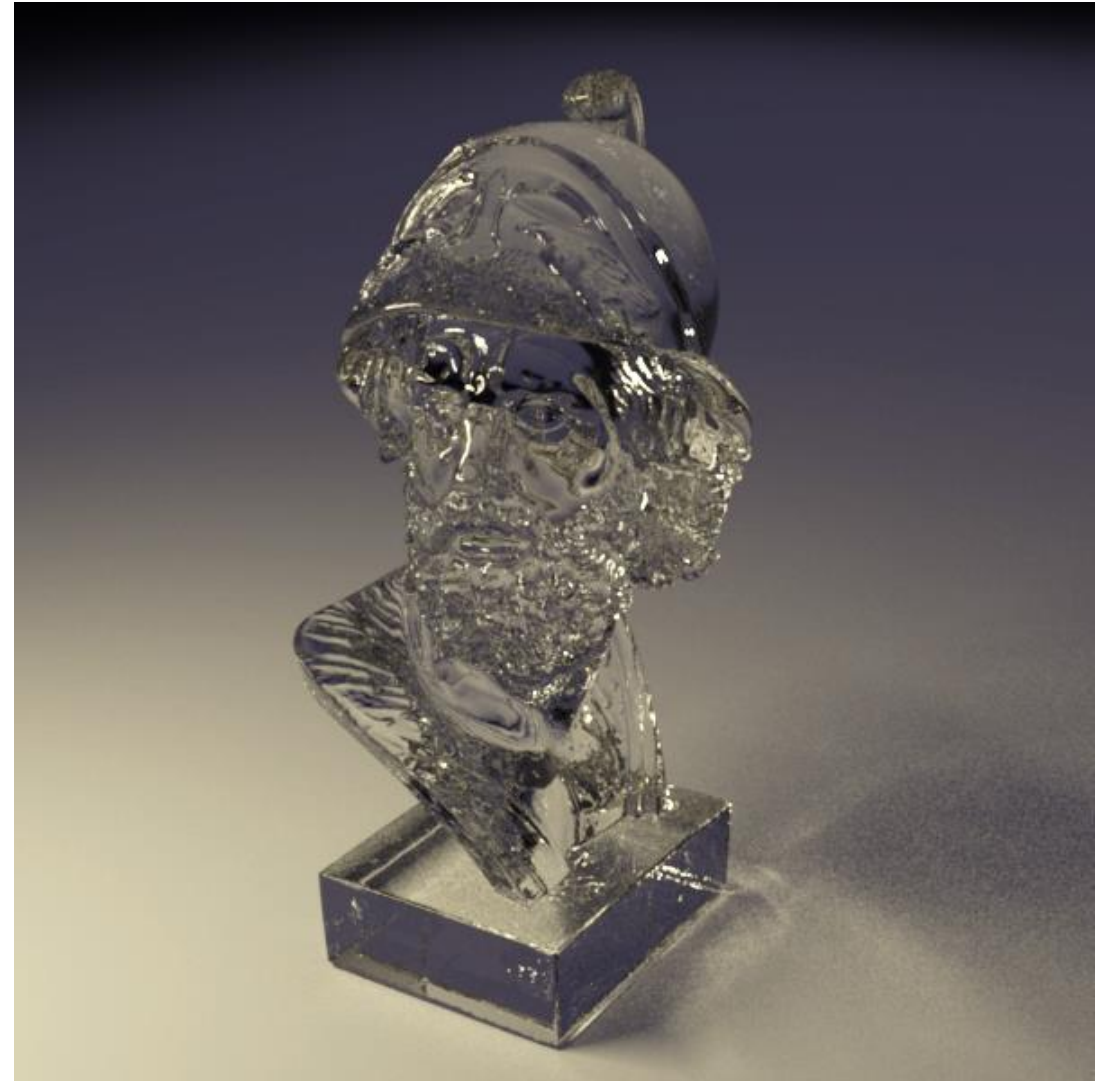
- **pdf(ω)**: always return 0



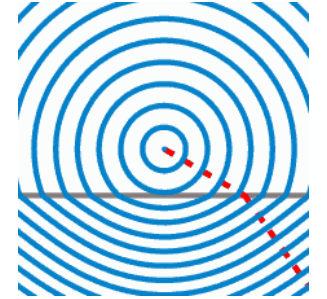


- Before, we assumed that the entire radiance is reflected (mirror)
- This is usually not the case
 - Some light is reflected on the surface
 - Some enters the new material (scattered, absorbed or **refracted**)
 - Meeting point of two different media is called **interface**
- When entering a different medium, light often changes direction
- Governed by the materials' index of refraction and *Snell's law*





- Based on the indices of refraction for the two materials
 - η_i for the medium that the light ray is currently in
 - η_t for the new medium into which light is transmitted

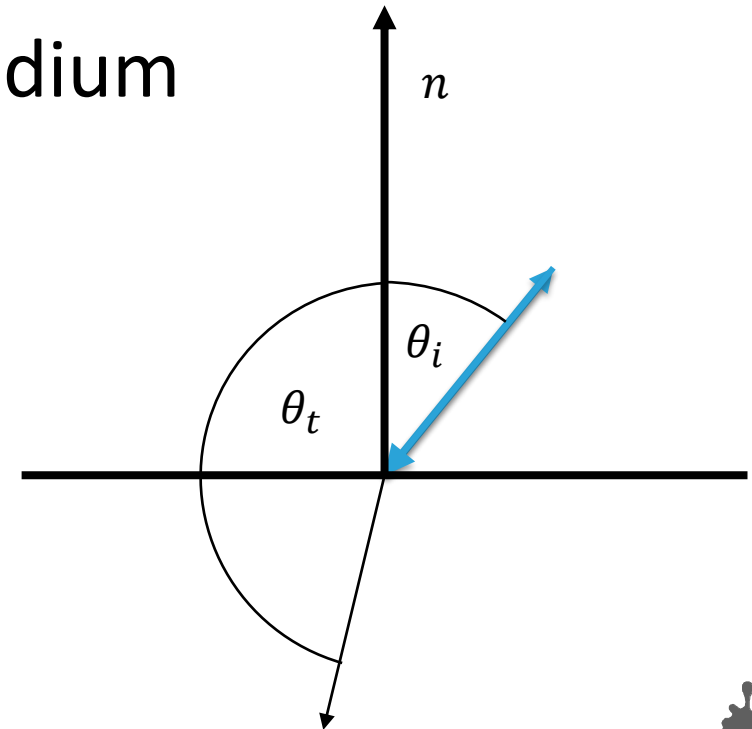


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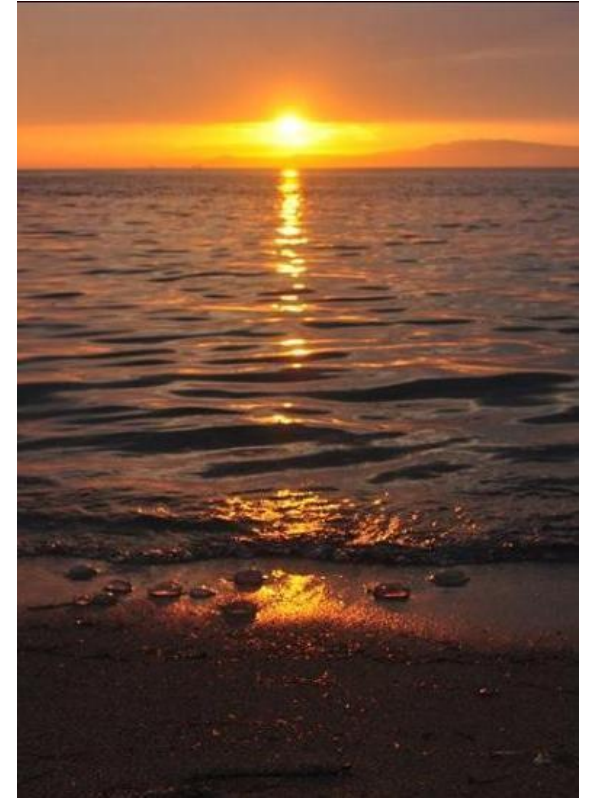
- Index of refraction: how fast light travels in medium
- Snell's law, essentially:

$$\eta_i \sin \theta_i = \eta_t \sin \theta_t$$

- Given η_i , θ_i and η_t , we can easily solve for θ_t



- How much of the light do we reflect?
- Not constant, but actually depends on the θ_i
- The larger θ_i , the better the chance for reflection
- If $\eta_i > \eta_t$, if incident light exceeds a certain θ_i , all light may be reflected (*total internal reflection*)

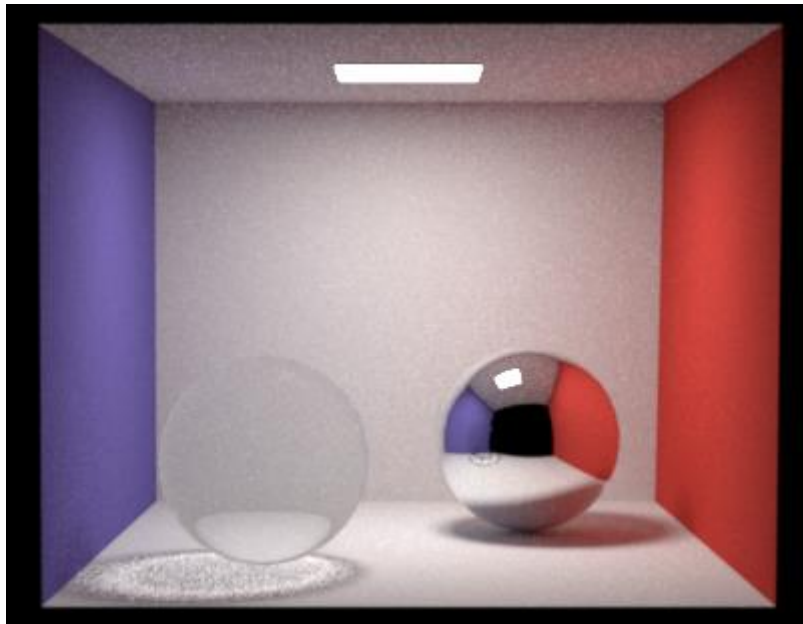


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Wikipedia, "Fresnel equations"

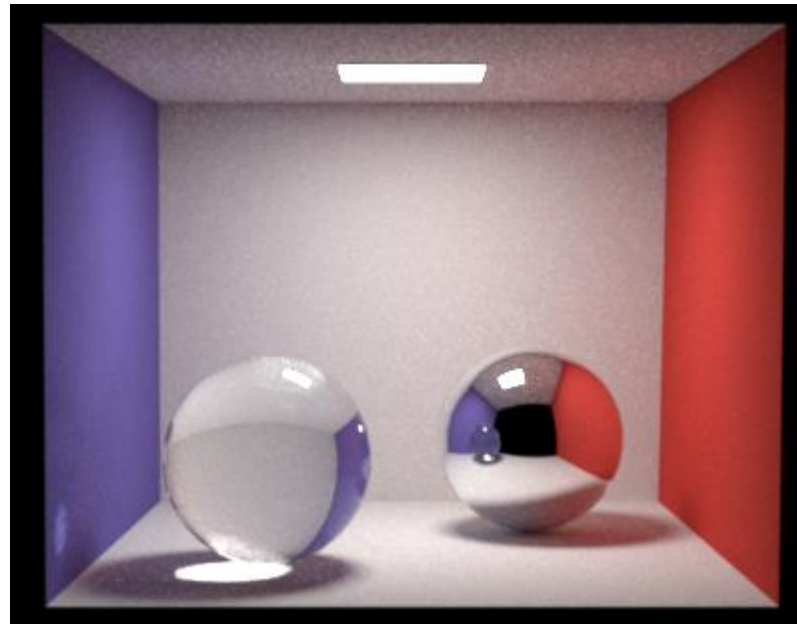


- Should be handled differently, depending on the materials involved
- Distinguish how material responds to energy transported by light
- We usually consider three major groups:
 - Dielectrics conduct electricity poorly (glass, air...)
 - Conductors (*metals*, reflect a lot, transmitted light quickly absorbed)
 - Semiconductors (complex, but also rare – we can ignore them)
- Remember, we are aiming for **dielectrics** today

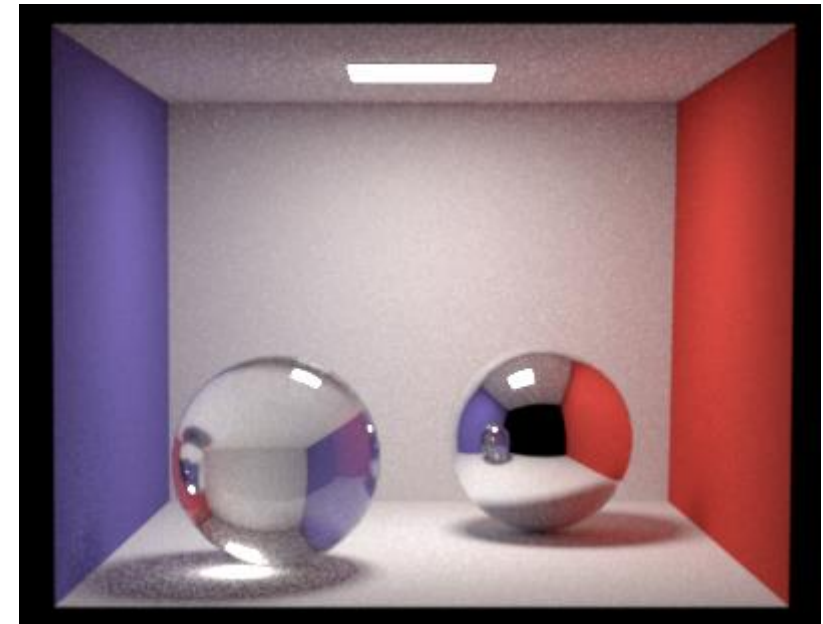




$\eta_t = 1.025$ (liquid helium)



$\eta_t = 1.5$ (glass)



$\eta_t = 2.5$ (diamond)

- Gases: 1 – 1.0005 (no-man's land from 1.05 to 1.25)
- Liquids: 1.3 (water) – 1.5 (olive oil)
- Solids: 1.3 (ice) – 2.5 (diamond)



- Defined for parallel and perpendicular polarized light (r_{\parallel} and r_{\perp}):

$$r_{\parallel} = \frac{\eta_t \cos \theta_i - \eta_i \cos \theta_t}{\eta_t \cos \theta_i + \eta_i \cos \theta_t}, r_{\perp} = \frac{\eta_i \cos \theta_i - \eta_t \cos \theta_t}{\eta_i \cos \theta_i + \eta_t \cos \theta_t}$$

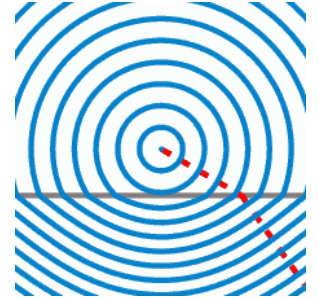
- Amount of **reflected** light (unpolarized light, average of squares):

$$F_r = \frac{1}{2} (r_{\parallel}^2 + r_{\perp}^2)$$

- Amount of **refracted** light (conservation of energy): $1 - F_r$



- Refracted light usually changes direction in new medium
 - Remember that we work with radiance: $d\Phi = L_i dA_{\perp} d\omega$
 - Refracted light changes direction \rightarrow influences radiance!
 - Relate incoming to refracted light:



Public domain, [Oleg Alexandrov](#),
Snell's law wavefronts, Wikipedia,
"Snell's law"

$$L_o \cos \theta_o dA \sin \theta_o d\theta_o d\phi_o = (1 - F_r) L_i \cos \theta_i dA \sin \theta_i d\theta_i d\phi_i$$

- Differentiating Snell's law w.r.t. θ , we get:

$$\eta_o \cos \theta_o d\theta_o = \eta_i \cos \theta_i d\theta_i \rightarrow \frac{\cos \theta_o d\theta_o}{\cos \theta_i d\theta_i} = \frac{\eta_i}{\eta_o}$$



- Substituting, we get:

$$L_o \eta_i^2 d\phi_o = (1 - F_r) L_i \eta_o^2 d\phi_i \rightarrow L_o = (1 - F_r) \frac{\eta_o^2}{\eta_i^2} L_i$$

- We have all the required information for the specular BTDF!
 - Use $T(\omega, n)$ to compute direction of ω when refracted at interface
 - Like specular BRDF, light only goes in a single direction
 - Can reuse BRDF $\delta(\omega)$ and normalization (similar implementation!)

$$f_r(x, \omega_i \rightarrow \omega_o) = \frac{\eta_o^2}{\eta_i^2} (1 - F_r) \frac{\delta(\omega_i - T(\omega_o, n))}{|\cos \theta_i|}$$



- When light refracts into a material with a higher η , the energy is compressed into a smaller set of angles
- For the BTDF, $f_r(x, \omega_i \rightarrow \omega_o) = f_r(x, \omega_o \rightarrow \omega_i)$ is not guaranteed
- No reciprocity, but $\eta_i^2 f_r(x, \omega_i \rightarrow \omega_o) = \eta_o^2 f_r(x, \omega_o \rightarrow \omega_i)$ holds!
- If you follow a **view ray**, do the same computations as above, just:
 - Make sure you choose η_i for medium ray comes from
 - Make sure you choose η_t for medium ray goes to



- Just continue one path, use Fresnel to decide \rightarrow reflect or refract?
- View ray behaves exactly like **incident light** in the above equations
- You may find it easier to flip the normal if light **exits** a medium
 - Light that enters e.g. a glass body must also exit at some point
 - I.e., the incoming light ray is not in same hemisphere as n
 - Consistent with using η_i and η_t for current/new medium
- Solving for θ_t , you may get “ $\sin \theta_t > 1$ ” \rightarrow **total internal reflection**

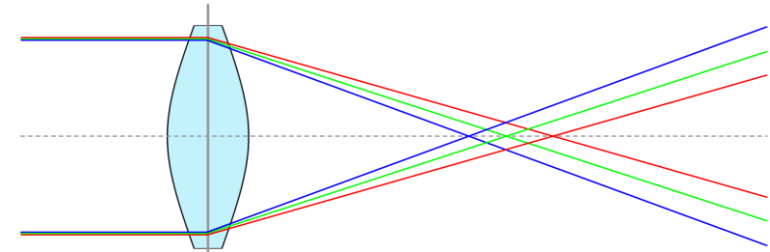


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- Physically speaking, the change in direction is wavelength-dependent
- For proper simulation, would have to at least bend R/G/B differently
- Would spawn two additional rays!
- Can of course be done, but is often ignored (tiny effect on most images)



Public domain, [Andreas 06](#), Chromatic aberration convex, Wikipedia, "Chromatic aberration"



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- Assign a letter to every interaction of a light path from light to eye
 - L – light
 - D – diffuse surface
 - S – specular surface
 - E – eye
- Use regex to describe specific (e.g., very challenging) path types
 - LE: direct path from light to eye
 - L(D|S)*E: any path from light to eye
 - LDS+E: a path with one diffuse bounce, followed by specular bounces



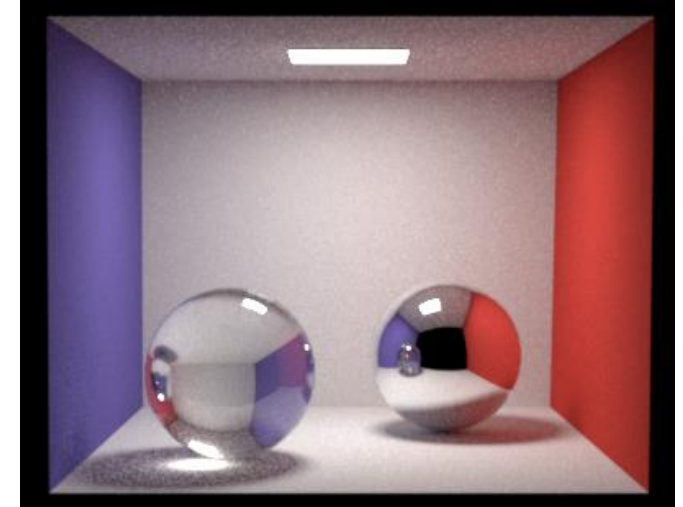
- General: focused light from interacting with curved, specular surface



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[CC BY-SA 4.0](#), Markus Selmke, Computer rendering of a wine glass caustic, Wikipedia, "Caustic (optics)"



- For us, who are concerned with rendering and path tracing: LS+DE
- Usually challenging to render (takes extremely long to converge)



- [SIGGRAPH University - Introduction to "Physically Based Shading in Theory and Practice"](#) by Naty Hoffman (!!!)
- [SIGGRAPH University - Recent Advances in Physically Based Shading](#) by Naty Hoffman
(advanced, in the same video there are also some other talks)



- *Material for Dielectrics largely based on “Physically Based Rendering” book, chapter 8: Reflection Models*
- [1] [Physically Based Rendering](#) (course book, chapters 8 and 9 for materials, chapter 11 for volume rendering)
- [2] [Background: Physics and Math of Shading](#) by Naty Hoffman
- [3] Wojciech Jarosz, “Efficient Monte Carlo Methods for Light Transport in Scattering Media”, PhD Thesis, <https://cs.dartmouth.edu/~wjarosz/publications/dissertation/>
- [4] [Production Volume Rendering](#) (SIGGRAPH 2017 Course)
- [5] [Monte Carlo methods for physically based volume rendering](#) (SIGGRAPH 2018 Course)

